One dimensional bound state Masatsugu Sei Suzuki Department of Physics, SUNY at Binghamton (Date: October 10, 2011)

1 One dimensional bound state

As a simple example of the calculation of discrete energy levels of a particle (with mass m) in quantum mechanics, we consider the one dimensional motion of a particle in the presence of a square-well potential barrier (width 2a and a depth V_0) as shown below.

V(x) = 0 for |x| > a, and $-V_0$ for -a < x < a.

If the energy of the particle E is negative, the particle is confined and in a bound state. Here we discuss the energy eigenvalues and the eigenfunctions for the bound states from the solution of the Schrödinger equation.



Fig.8 One dimensional square well potential of width 2a and depth V_0 .

(a) The parity of the wave function

When potential is an even function (symmetric with respect to x), the wave function should have even parity or odd parity.

((Proof))

$$[\hat{\pi},\hat{H}]=0$$

 $\hat{\pi}$ is the parity operator.

- $\hat{\pi}^2 = 1$ $\hat{\pi}^+ = \hat{\pi} = \hat{\pi}^{-1}.$
- $\hat{\pi}\hat{x}\hat{\pi} = -\hat{x} \,. \qquad \hat{\pi}\hat{p}\,\hat{\pi} = -\hat{p} \,.$

 \hat{H} is the Hamiltonian.

$$\hat{H} = \frac{\hat{p}^2}{2m} + V(\hat{x}),$$

and

$$\hat{\pi}\hat{H}\hat{\pi} = \hat{\pi}[\frac{\hat{p}^{2}}{2m} + V(\hat{x})]\hat{\pi}$$

$$= \frac{1}{2m}(\hat{\pi}\hat{p}\hat{\pi})^{2} + V(\hat{\pi}\hat{x}\hat{\pi})$$

$$= \frac{1}{2m}(-\hat{p})^{2} + V(-\hat{x})$$

$$= \frac{1}{2m}\hat{p}^{2} + V(\hat{x})$$

since $V(-\hat{x}) = V(\hat{x})$. Then we have a simultaneous eigenket:

$$\hat{H}|\psi\rangle = E|\psi\rangle$$
, and $\hat{\pi}|\psi\rangle = \lambda|\psi\rangle$.

Since $\hat{\pi}^2 = 1$,

$$\hat{\pi}^2 |\psi\rangle = \lambda \hat{\pi} |\psi\rangle = \lambda^2 |\psi\rangle = |\psi\rangle.$$

Thus we have $\lambda = \pm 1$.

or

$$\begin{aligned} \hat{\pi} |\psi\rangle &= \pm |\psi\rangle, \\ \langle x |\hat{\pi} |\psi\rangle &= \pm \langle x |\psi\rangle. \end{aligned}$$

Since

$$\hat{\pi}|x\rangle = |-x\rangle$$
, or $\langle x|\hat{\pi}^+ = \langle x|\hat{\pi} = \langle -x|$

we have

$$\langle -x|\psi\rangle = \pm \langle x|\psi\rangle,$$

$$\psi(-x) = \pm \psi(x) \, .$$

(b) Wavefunctions

In the Regions I, II, and III, the Schrödinger equation takes the form

$$\frac{d^2}{dx^2}\psi(x) - \kappa^2\psi(x) = 0$$
 outside the well.
$$\frac{d^2}{dx^2}\psi(x) + k^2\psi(x) = 0$$
 inside the well.

Here we define

$$\kappa^{2} = \frac{2m}{\hbar^{2}} |E|, \qquad k^{2} = \frac{2m}{\hbar^{2}} (V_{0} - |E|).$$

Here we introduce parameters (β and σ) for convenience,

$$\kappa^{2} = \frac{2m}{\hbar^{2}} |E| = \frac{2mV_{0}}{\hbar^{2}} \frac{|E|}{V_{0}} = \frac{2mV_{0}a^{2}}{\hbar^{2}} \frac{1}{a^{2}} \frac{|E|}{V_{0}} = \frac{\beta^{2}}{a^{2}} \varepsilon,$$

or

$$\kappa^2 = \frac{\beta^2}{a^2} \varepsilon \,,$$

and

$$k^{2} = \frac{2m}{\hbar^{2}} (V_{0} - |E|) = \frac{2mV_{0}}{\hbar^{2}} (1 - \frac{|E|}{V_{0}}) = \frac{1}{a^{2}} \beta^{2} (1 - \varepsilon) ,$$

where

$$\varepsilon = \frac{|E|}{V_0}$$
, and $\beta = \sqrt{\frac{2mV_0a^2}{\hbar^2}}$.

We note that

$$k^2 + \kappa^2 = \frac{\beta^2}{a^2},$$

$$\xi^2 + \eta^2 = \beta^2,$$

where $ka = \xi$ and $ka = \eta$. The energy ε is given by

$$\varepsilon = \frac{\eta^2}{\beta^2} = 1 - \frac{\xi^2}{\beta^2} \,.$$

The stationary solution of the three regions are given by

$$\begin{split} \varphi_I(x) &= A e^{\kappa x} \,, \\ \varphi_{II}(x) &= B_1 e^{ikx} + B_2 e^{-ikx} \,, \\ \varphi_{III}(x) &= C e^{-\kappa x} \,. \end{split}$$

(i) The wave function with even parity

$$A = C,$$

$$B_1 = B_2 \equiv \frac{B}{2}.$$

The wavefunctions can be described by

$$\varphi_{I}(x) = Ae^{\kappa x},$$
$$\varphi_{II}(x) = B\cos(kx),$$
$$\varphi_{III}(x) = Ae^{-\kappa x}.$$

The derivatives are obtained by

$$\frac{d\varphi_{I}(x)}{dx} = A\kappa e^{\kappa x},$$
$$\frac{d\varphi_{II}(x)}{dx} = -Bk\sin(kx),$$
$$\frac{d\varphi_{III}(x)}{dx} = -A\kappa e^{-\kappa x}.$$

At x = a, $\varphi(x)$ and $\frac{d\varphi(x)}{dx}$ are continuous. Then we have $Ae^{-\kappa a} - B\cos(ka) = 0$, $-A\kappa e^{-\kappa a} + Bk\sin(ka) = 0$,

or

$$MX=0,$$

where

$$M = \begin{pmatrix} e^{-\kappa a} & -\cos(ka) \\ -\kappa e^{-\kappa a} & k\sin(ka) \end{pmatrix}, \qquad \qquad X = \begin{pmatrix} A \\ B \end{pmatrix}.$$

The condition det*M*=0 leads to

$$k\sin(ka)e^{-\kappa a}=\kappa e^{-\kappa a}\cos(ka),$$

or

$$\tan(ka) = \frac{\kappa}{k}$$
 for the even parity,

or

$$\kappa a = ka \tan(ka)$$
 for the even parity.

or

$$\eta = \xi \tan \xi$$
.

The constants A, B, and C are given by

$$A = C = Be^{\kappa a} \cos(ka).$$

The condition of the normalization leads to the value of B.

(ii) The wave function with odd parity

A = -C,

$$B_1 = -B_2 \equiv \frac{B}{2i} \, .$$

The wavefunctions are given by

$$\varphi_{II}(x) = -Ae^{\kappa x},$$
$$\varphi_{II}(x) = B\sin(kx),$$
$$\varphi_{III}(x) = Ae^{-\kappa x}.$$

The derivatives are obtained as

$$\frac{d\varphi_{I}(x)}{dx} = -A\kappa e^{\kappa x},$$
$$\frac{d\varphi_{II}(x)}{dx} = Bk\cos(kx),$$
$$\frac{d\varphi_{III}(x)}{dx} = -A\kappa e^{-\kappa x}.$$

At x = a, $\varphi(x)$ and $\frac{d\varphi(x)}{dx}$ are continuous. Then we have

$$-Ae^{-\kappa a} + B\sin(ka) = 0,$$
$$-A\kappa e^{-\kappa a} - Bk\cos(ka) = 0$$

or

$$MX=0,$$

where

$$M = \begin{pmatrix} -e^{-\kappa a} & \sin(ka) \\ -\kappa e^{-\kappa a} & -k\cos(\frac{ka}{2}) \end{pmatrix}, \qquad X = \begin{pmatrix} A \\ B \end{pmatrix}.$$

,

The condition det*M*=0 leads to

$$k\cos(ka)e^{-\kappa a} = -\kappa e^{-\kappa a}\sin(ka),$$

 $\kappa a = -ka \cot(ka)$ for the odd parity,

or

or

$$\eta = -\xi \cot \xi \ .$$

We solve this eigenvalue problem using the Mathematica. The result is as follows.



Fig.9 Graphical solution. One solution with even parity for $0 < \beta < \pi/2$. One solution with even parity and one solution with odd parity for $\pi/2 < \beta < \pi$. Two solutions with even parity and one solution with odd parity for $\pi < \beta < 3\pi/2$. Two solutions with even parity and two solutions with odd parity for $3\pi/2 < \beta < 2\pi$. $\eta = \xi \tan \xi$ for the even parity (red lines). $\eta = -\xi \cot \xi$ for the odd parity (blue lines). The circles are denoted by $\xi^2 + \eta^2 = \beta^2$. The parameter β is changed as $\beta = 1, 2, 3, 4$, and 5. $\varepsilon = \frac{|E|}{V_0} = \frac{\eta^2}{\beta^2} = 1 - \frac{\xi^2}{\beta^2}$. $\xi = ka$ and $\eta = \kappa a$.

The normalized wavefunction for the even parity and odd parity are given by

$$\psi \mathbf{eI} = \frac{\mathbf{e}^{\eta + \mathbf{x} \eta} \operatorname{Cos}[\xi]}{\sqrt{1 + \frac{\cos[\xi]^2}{\eta} + \frac{\sin[2\xi]}{2\xi}}}; \quad \psi \mathbf{eII} = \frac{\operatorname{Cos}[\mathbf{x} \xi]}{\sqrt{1 + \frac{\cos[\xi]^2}{\eta} + \frac{\sin[2\xi]}{2\xi}}};$$
$$\psi \mathbf{eIII} = \frac{\mathbf{e}^{\eta - \mathbf{x} \eta} \operatorname{Cos}[\xi]}{\sqrt{1 + \frac{\cos[\xi]^2}{\eta} + \frac{\sin[2\xi]}{2\xi}}};$$
$$\psi \mathbf{oII} = -\frac{\mathbf{e}^{\eta + \mathbf{x} \eta} \operatorname{Sin}[\xi]}{\sqrt{1 + \frac{\sin[\xi]^2}{\eta} - \frac{\sin[2\xi]}{2\xi}}}; \quad \psi \mathbf{oII} = \frac{\operatorname{Sin}[\mathbf{x} \xi]}{\sqrt{1 + \frac{\sin[\xi]^2}{\eta} - \frac{\sin[2\xi]}{2\xi}}};$$
$$\psi \mathbf{oIII} = \frac{\mathbf{e}^{\eta - \mathbf{x} \eta} \operatorname{Sin}[\xi]}{\sqrt{1 + \frac{\sin[\xi]^2}{\eta} - \frac{\sin[2\xi]}{2\xi}}};$$

for the regions I, II, and III, where ψ_e is the wavefunction with the even parity and ψ_o is the wavefunction with the odd parity.

$\beta = 1$			
$\xi_{11} = 0.739085$	$\eta_{11} = 0.673612$	$\varepsilon_{II} = 0.453753$	3 even
$\overline{\beta}=2$			
$\xi_{21} = 1.02987$ $\xi_{22} = 1.89549$	$\eta_{21} = 1.71446$ $\eta_{22} = 0.638045$	$\varepsilon_{21} = 0.734844$ $\varepsilon_{22} = 0.101775$	even odd
$\overline{\beta}=3$			
$\xi_{31} = 1.17012$ $\xi_{32} = 2.27886$	$\eta_{31} = 2.76239$ $\eta_{32} = 1.9511$	$\varepsilon_{31} = 0.847869$ $\varepsilon_{32} = 0.422976$	even odd
$\beta = 4$			
$\xi_{41} = 1.25235$ $\xi_{42} = 2.47458$ $\xi_{43} = 3.5953$	$\eta_{41} = 3.7989$ $\eta_{42} = 3.14269$ $\eta_{43} = 1.75322$	$\varepsilon_{41} = 0.901976$ $\varepsilon_{42} = 0.617279$ $\varepsilon_{43} = 0.192111$	even odd even

$\xi_{51} = 1.30644$	$\eta_{51} = 4.8263,$	$\varepsilon_{51} = 0.931729$	even
$\xi_{52} = 2.59574$	$\eta_{52} = 4.27342,$	$\varepsilon_{52} = 0.730486$	odd
$\xi_{53} = 3.83747$	$\eta_{53} = 3.20528,$	$\varepsilon_{53} = 0.410954$	even
$\xi_{54} = = 4.9063$	$\eta_{54} = .963467,$	$\varepsilon_{54} = 0.0371307$	odd

 $\beta = 5$



Fig.10 Square well potential V(x) of width 2a and depth V_0 . $\beta = 1$ and the corresponding wavefunction $\psi(x)$ which is normalized. There is one bound state (even parity) (- $\varepsilon_{11} = -0.45735$), where $\varepsilon = |E|/V_0$.



Fig.11 $\beta = 2$. There are two bound states. (i) The bound state (denoted by red) with even parity (- $\varepsilon_{21} = -0.734844$). (ii) The bound state (denoted by blue) with odd parity (- $\varepsilon_{22} = -0.101775$).



Fig.12 $\beta = 3$. There are two bound states. (i) The bound state (denoted by red) with even parity ($-\varepsilon_{31} = -0.847869$). (ii) The bound state (denoted by blue) with odd parity ($-\varepsilon_{32} = -0.422976$).



Fig.13 $\beta = 4$. There are three bound states. (i) The bound state (denoted by red) with even parity (- $\varepsilon_{41} = -0.901976$). (ii) The bound state (denoted by blue)

with odd parity (- ε_{42} = -0.617279). (iii) The bound state (denoted by red) with even parity (- ε_{43} = -0.192111).



Fig.14 $\beta = 5$. There are four bound states. (i) The bound state (denoted by red) with even parity ($-\varepsilon_{51} = -0.931729$). (ii) The bound state (denoted by blue) with odd parity ($-\varepsilon_{52} = -0.730486$). (iii) The bound state (denoted by red) with even parity ($-\varepsilon_{53} = -0.410954$). (iv) The bound state (denoted by blue) with odd parity ($-\varepsilon_{54} = -0.0371307$).

REFERENCES

- 1. L.I. Schiff, Quantum Mechanics (McGraw-Hill, New York, 1955).
- 2. E. Merzbacher, Quantum Mechanics Third edition (John Wiley and Sons, New York, 1998.

II. Bound states

Bound state-II



We suppose that the potential energy is given by

 $V(x) = a |x| \qquad (a>0)$

The Schrödinger equation is given by

$$\frac{d^2\psi}{dx^2} + \frac{2m}{\hbar^2}(E - a|x|)\psi = 0$$

where E is the energy of a particle with a mass m.

Since the potential is a even function of x. the wave function should be either an even function or a odd function of x.

The boundary condition for the wave function with the odd parity;

 $\psi(x) = 0 \qquad \text{at } x = 0.$

The boundary condition for the wave function with even parity,

$\frac{d\psi}{dx} = 0$	at $x = 0$.
n = 0	even parity
n=0.	even parity
n = 1	odd parity
n = 2	even parity
<i>n</i> = 3	odd parity
n = 4	even parity



We use the dimensionless parameters;

$$\varepsilon = \frac{E}{\left(\hbar^2 a^2 / m\right)^{1/3}},$$
$$z = \frac{x}{\left(\hbar^2 / ma\right)^{1/3}}.$$

We note that

$$\frac{d\psi}{dx} = \frac{dz}{dx}\frac{d\psi}{dz} = \frac{1}{\left(\frac{\hbar^2}{ma}\right)^{1/3}}\frac{d\psi}{dz}$$
$$\frac{d^2\psi}{dx^2} = \frac{dz}{dx}\frac{d}{dz}\frac{d\psi}{dx} = \frac{1}{\left(\frac{\hbar^2}{ma}\right)^{2/3}}\frac{d^2\psi}{dz^2}$$

Then we get a differential equation

$$\frac{d^2\psi}{dz^2} + \frac{2m}{\hbar^2} (\hbar^2 / ma)^{2/3} [(\hbar^2 a^2 / m)^{1/3} \varepsilon - a |z| (\hbar^2 / ma)^{1/3}] \psi = 0$$

or

$$\frac{d^2\psi}{dz^2} + 2(\varepsilon - |z|)\psi = 0$$

Since the wave function either an even function or odd function, we consider the case of z>0.

$$\frac{d^2\psi}{dz^2} + 2(\varepsilon - z)\psi = 0$$

Here we put $y = z - \varepsilon$.

$$\frac{d^2\psi(y)}{dy^2} - 2y\psi(y) = 0$$

((Mathematica))

```
Clear["Global`*"];
g1 = D[y[x], {x, 2}] - 2 x y[x] == 0;
eq1 = DSolve[g1, y[x], x]
{{y[x] → AiryAi[2<sup>1/3</sup> x] C[1] + AiryBi[2<sup>1/3</sup> x] C[2]}}
y1[x_] = y[x] /. eq1[[1]] /. {C[2] → 0, C[1] → 1}
AiryAi[2<sup>1/3</sup> x]
```

The solution of this differential equation is obtained as

$$\psi(y) = cA_i(2^{1/3}y)$$

where A_i is an Airy function and c is a constant to be determined from the normalization. Note that the second solution B_i is not a solution in this case since the function diverges when $x \rightarrow \infty$

(a) The case of even function (the even parity)

The boundary condition:
$$\frac{d\psi(z)}{dz}$$
 at $z = 0$.

which means that

$$\frac{d\psi(z)}{dz} = \frac{d\psi(y)}{dy} = 2^{1/3} A_i'(2^{1/3} y) = 0 \qquad \text{at } y = -\varepsilon.$$



From this we get an energy eigenvalue for the wave functions with the even parity. The points with Green are located at $x = -\varepsilon_0, -\varepsilon_2, -\varepsilon_4, -\varepsilon_6,...$

 $\varepsilon_0 = 0.808617$

$$\psi_0(z) = 1.468A_i(2^{1/3}(z-\varepsilon_0))$$

 $\varepsilon_2 = 2.5781$

$$\psi_2(z) = 1.0510A_i(2^{1/3}(z-\varepsilon_2))$$

 $\varepsilon_4 = 3.82572$

$$\psi_4(z) = 1.0510 A_i (2^{1/3} (z - \varepsilon_4))$$

 $\varepsilon_6 = 4.89182$ $\varepsilon_8 = 5.8513$ $\varepsilon_{10} = 6.73732$ $\varepsilon_{12} = 7.56829$ $\varepsilon_{14} = 8.35581$

(i)



The wave function is normalized.



(b) The case of odd function:

The boundary condition: $\psi(z)$ at z = 0.

which means that

$$\psi(y) = 0$$
 at $y = -\varepsilon$.

From this we get an energy eigenvalue for the wave functions with the odd parity. The points with Green are located at $x = -\varepsilon_1, -\varepsilon_3, -\varepsilon_5, -\varepsilon_7, \dots$



 $\varepsilon_1 = 1.85576$

$$\psi_1(z) = 1.1319A_i(2^{1/3}(z-\varepsilon_1))$$

 $\varepsilon_3 = 3.24461$

$$\psi_3(z) = 0.988282A_i(2^{1/3}(z-\varepsilon_3))$$

 $\varepsilon_5 = 4.38167$ $\varepsilon_7 = 5.38661$ $\varepsilon_9 = 6.30526$ $\varepsilon_{11} = 7.16128$ $\varepsilon_{13} = 7.96889$

(i)

 $\varepsilon_1 = 1.85576$ (*n* = 1) odd parity







The plot of wave functions (n = 0, 1, 2, 3). The blue lines show the energy levels.