

**Polarizer and the photon polarization**  
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For the Stern-Gerlach experiment with spin 1/2, after the measurement of the spin direction  $S_z$ , the state of the system collapses into one of two states,  $|+z\rangle$  and  $|-z\rangle$ . The situation is rather different for the light polarization, the state of the system collapses into one state depending on the direction of polarizer such as  $x$ -axis polarizer; the eigenket  $|x\rangle$  for the  $x$ -polarizer, and the eigenket  $|y\rangle$  for the  $y$ -polarizer. The final state is given by the projection operator which is applied to the initial state  $|\psi\rangle$ . Here we discuss the role of the projection operator for the photon polarization.

**1. Stern-Gerlach experiment for the spin 1/2 (as an example)**

First we consider the Stern-Gerlach experiment ( $SG_z$ ). We have two eigenkets of the spin operator  $\hat{S}_z$ ,  $|+z\rangle$  and  $|-z\rangle$ .

$$\hat{S}_z|+z\rangle = \frac{\hbar}{2}|+z\rangle, \quad \hat{S}_z|-z\rangle = -\frac{\hbar}{2}|-z\rangle$$

The projection operators are defined by

$$\hat{P}_{+z} = |+z\rangle\langle +z|, \quad \hat{P}_{-z} = |-z\rangle\langle -z|$$

The spin operator  $\hat{S}_z$  can be expressed by

$$\hat{S}_z = \frac{\hbar}{2}(|+z\rangle\langle +z| - |-z\rangle\langle -z|) = \frac{\hbar}{2}(\hat{P}_{+z} - \hat{P}_{-z})$$

(a) **The measurements: eigenvalue problem**

The eigenkets  $|+z\rangle$  and  $|-z\rangle$  are determined from the eigenvalue problem. After the measurements the system collapses

(b) **Projection operator**

When the initial state of the system is given by  $|\psi\rangle$ , the final states after the  $SG_z$  are

$$\hat{P}_{+z}|\psi\rangle, \quad \hat{P}_{-z}|\psi\rangle$$

using the projection operators.

(c) **The probability**

The probability of finding the system in the state  $|+z\rangle$  is

$$\langle +z | \hat{P}_{+z} | \psi \rangle = \langle +z | \psi \rangle$$

The probability of finding the system in the state  $|-z\rangle$  is

$$\langle -z | \hat{P}_{-z} | \psi \rangle = \langle -z | \psi \rangle$$

**2. Photon polarization**

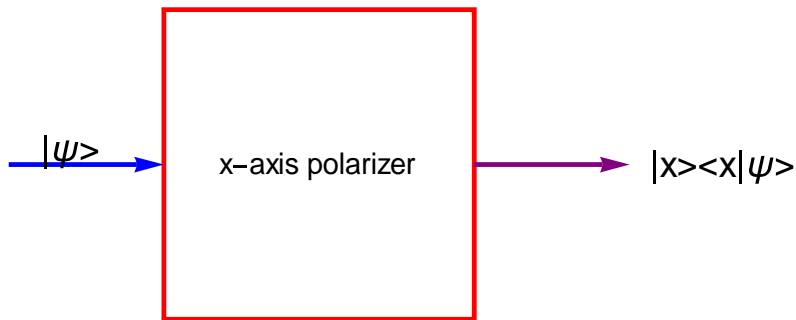
(a) **Projection operator**

The final state of the light after passing the  $x$ -axis polarizer is given by

$$\hat{P}_x|\psi\rangle = |x\rangle\langle x|\psi\rangle.$$

The probability of finding the system in the state  $|x\rangle$  is

$$|\langle x | \hat{P}_x | \psi \rangle|^2 = |\langle x | \psi \rangle|^2.$$



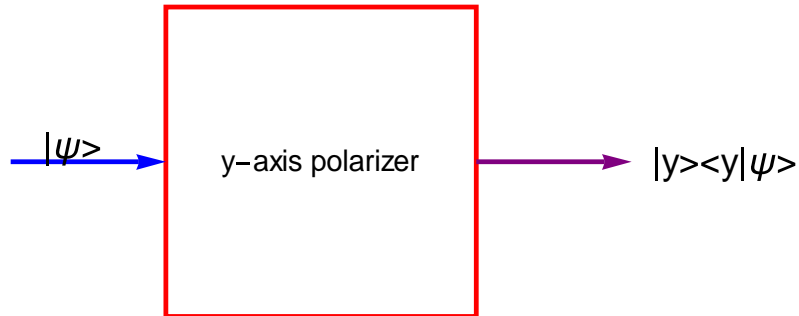
(b) **The use of y-axis polarizer**

The final state of the light after passing the  $y$ -axis polarizer is given by

$$\hat{P}_y|\psi\rangle = |y\rangle\langle y|\psi\rangle.$$

The probability of finding the system in the state  $|y\rangle$  is

$$|\langle y|\hat{P}_y|\psi\rangle|^2 = |\langle y|\psi\rangle|^2.$$



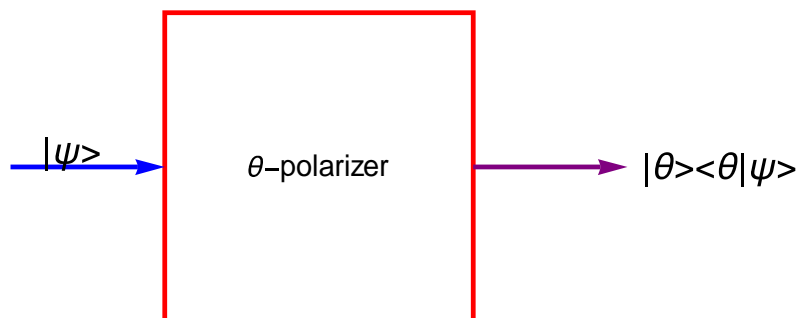
(c) **Th polarizer with the angle  $\theta$**

The final state of the light after passing the angle  $\theta$  polarizer is given by

$$\hat{P}_\theta|\psi\rangle = |\theta\rangle\langle\theta|\psi\rangle.$$

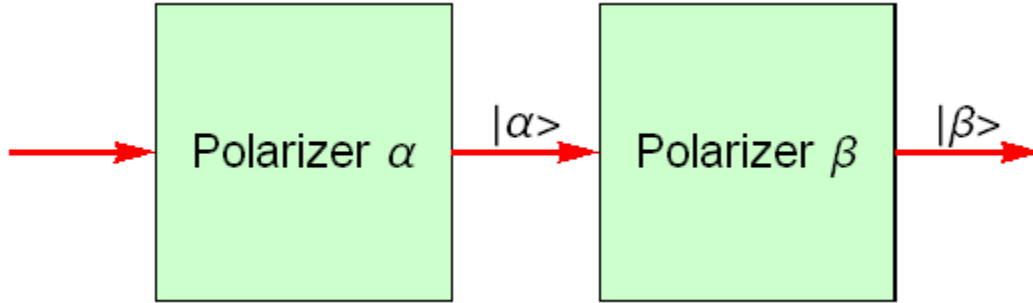
The probability of finding the system in the state  $|\theta\rangle$  is

$$|\langle\theta|\hat{P}_\theta|\psi\rangle|^2 = |\langle\theta|\psi\rangle|^2.$$



### 3. Examples-1

Suppose we use two polarizers (angles  $\alpha$  and  $\beta$ ) in series;  $\alpha - \beta$



The projection operators for the polarizers  $\alpha$  and  $\beta$  are

$$\hat{P}_\alpha = |\alpha\rangle\langle\alpha| = \begin{pmatrix} \cos\alpha \\ \sin\alpha \end{pmatrix} \begin{pmatrix} \cos\alpha & \sin\alpha \end{pmatrix} = \begin{pmatrix} \cos^2\alpha & \sin\alpha\cos\alpha \\ \sin\alpha\cos\alpha & \sin^2\alpha \end{pmatrix}$$

and

$$\hat{P}_\beta = |\beta\rangle\langle\beta| = \begin{pmatrix} \cos\beta \\ \sin\beta \end{pmatrix} \begin{pmatrix} \cos\beta & \sin\beta \end{pmatrix} = \begin{pmatrix} \cos^2\beta & \sin\beta\cos\beta \\ \sin\beta\cos\beta & \sin^2\beta \end{pmatrix}$$

where

$$|\alpha\rangle = \begin{pmatrix} \cos\alpha \\ \sin\alpha \end{pmatrix}, \quad |\beta\rangle = \begin{pmatrix} \cos\beta \\ \sin\beta \end{pmatrix}$$

The initial state is given by  $|\psi\rangle$ . After passing through two polarizers  $\alpha$  and  $\beta$ , the final state is obtained as

$$\hat{P}_\beta \hat{P}_\alpha |\psi\rangle$$

The probability amplitude for the system in the final state  $|\beta\rangle$ , is given by

$$\langle\beta|\hat{P}_\beta\hat{P}_\alpha|\psi\rangle = \langle\beta|\alpha\rangle\langle\alpha|\psi\rangle$$

The corresponding probability is

$$P_{\alpha\beta} = |\langle\beta|\alpha\rangle|^2 |\langle\alpha|\psi\rangle|^2 = |\langle\alpha|\psi\rangle|^2 \cos^2(\beta - \alpha). \quad (\text{Malus' law})$$

where

$$\langle\beta|\alpha\rangle = (\cos\beta \quad \sin\beta) \begin{pmatrix} \cos\alpha \\ \sin\alpha \end{pmatrix} = \cos(\beta - \alpha)$$

#### 4. Example-II

Next we consider  $n$  polarizers (angles  $\alpha_1, \alpha_2, \alpha_3, \dots, \alpha_{n-1}, \alpha_n$ ) in series. The initial state is given by  $|\psi\rangle$ . After passing through  $n$  polarizers (angles  $\alpha_1, \alpha_2, \alpha_3, \dots, \alpha_{n-1}, \alpha_n$ ) in series, the final state can be obtained as

$$\hat{P}_{\alpha_n} \hat{P}_{\alpha_{n-1}} \dots \hat{P}_{\alpha_4} \hat{P}_{\alpha_3} \hat{P}_{\alpha_2} \hat{P}_{\alpha_1} |\psi\rangle$$

The probability amplitude for the system in the final state  $|\alpha_n\rangle$ , is given by

$$\langle\alpha_n | \hat{P}_{\alpha_n} \hat{P}_{\alpha_{n-1}} \dots \hat{P}_{\alpha_4} \hat{P}_{\alpha_3} \hat{P}_{\alpha_2} \hat{P}_{\alpha_1} |\psi\rangle = \langle\alpha_n | \alpha_{n-1}\rangle \dots \langle\alpha_3 | \alpha_2\rangle \langle\alpha_2 | \alpha_1\rangle \langle\alpha_1 | \psi\rangle$$

where

$$\hat{P}_{\alpha_k} = |\alpha_k\rangle\langle\alpha_k| = \begin{pmatrix} \cos^2 \alpha_k & \sin \alpha_k \cos \alpha_k \\ \sin \alpha_k \cos \alpha_k & \sin^2 \alpha_k \end{pmatrix}$$

The corresponding probability is

$$P = |\langle\alpha_1 | \psi\rangle|^2 \cos^2(\alpha_n - \alpha_{n-1}) \cos^2(\alpha_{n-1} - \alpha_{n-2}) \dots \cos^2(\alpha_3 - \alpha_2) \cos^2(\alpha_2 - \alpha_1).$$

#### REFERENCES

Richard P. Feynman and Albert R. Hibbs, *Quantum Mechanics and Path Integrals*, emended by Daniel F. Styer, Emended edition (Dover Publications, Inc. New York, 2010).