

Differential operators in Cartesian coordinates
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Here we discuss the differential operators for the Cartesian coordinates with the use of Mathematica. To this end, we use the problems of Goswami (quantum mechanics) and Binney and Skipper (quantum mechanics).

1. **Cartesian coordinates (Mathematica)**

Here we use the following notations in the Mathematica program for the spherical coordinates.

$$\text{Lap} \rightarrow \nabla^2$$

$$\text{Gra} \rightarrow \text{grad}, \quad \nabla$$

$$\text{Diva} \rightarrow \text{div}$$

$$\text{Curla} \rightarrow \text{curl} \quad \text{or} \quad \text{rot}$$

$$L \rightarrow \mathbf{L} = \mathbf{r} \times \mathbf{p}$$

$$L_x \rightarrow \mathbf{e}_x \cdot \mathbf{L}$$

$$L_y \rightarrow \mathbf{e}_y \cdot \mathbf{L}$$

$$L_z \rightarrow \mathbf{e}_z \cdot \mathbf{L}$$

$$LP \rightarrow L_x + iL_y$$

$$LM \rightarrow L_x - iL_y$$

$$Leq \rightarrow \mathbf{L}^2 = L_x^2 + L_y^2 + L_z^2$$

$$p \rightarrow \mathbf{p} = \frac{\hbar}{i} \nabla$$

$$p_x \rightarrow \mathbf{e}_x \cdot \mathbf{p} = \frac{\hbar}{i} \mathbf{e}_x \cdot \nabla$$

$$p_y \rightarrow \mathbf{e}_y \cdot \mathbf{p} = \frac{\hbar}{i} \mathbf{e}_y \cdot \nabla .$$

$$p_z \rightarrow \mathbf{e}_z \cdot \mathbf{p} = \frac{\hbar}{i} \mathbf{e}_z \cdot \nabla .$$

$$Peq \rightarrow \mathbf{p}^2 = p_x^2 + p_y^2 + p_z^2 .$$

$$p_r \rightarrow \frac{1}{2} \left(\frac{\mathbf{r} \cdot \mathbf{p}}{r} + \mathbf{p} \cdot \frac{\mathbf{r}}{r} \right)$$

$$\mathbf{r} \rightarrow x\mathbf{e}_x + y\mathbf{e}_y + z\mathbf{e}_z$$

$$R \rightarrow \sqrt{x^2 + y^2 + z^2}.$$

((Mathematica)) Differential operators in Cartesian coordinates

```

Clear["Global`"];
ux = {1, 0, 0};
uy = {0, 1, 0};
uz = {0, 0, 1};
r = {x, y, z};
R = Sqrt[r.r];
L := (-I h (Cross[r, Grad[#, {x, y, z}, "Cartesian"]]) &) // Simplify;
p = (-I h Grad[#, {x, y, z}, "Cartesian"] &) // Simplify;
Lx := (ux.L[#]) &;
Ly := (uy.L[#]) &;
Lz := (uz.L[#]) &;
px := (ux.p[#]) &;
py := (uy.p[#]) &;
pz := (uz.p[#]) &;
Lap := Laplacian[#, {x, y, z}, "Cartesian"] &;
Gra := Grad[#, {x, y, z}, "Cartesian"] &;
Cur := Curl[#, {x, y, z}, "Cartesian"] &;
Diva := Div[#, {x, y, z}, "Cartesian"] &;
prq := (-I h) 1/2 (r.R.Gra[#] + Diva[r/R #]) &;
Lsq := (Nest[Lx, #, 2] + Nest[Ly, #, 2] + Nest[Lz, #, 2]) &;

```

2. Commutation relation (I)

Show that

$$[\hat{L}_i, \hat{x}_j] = i\hbar \sum_k \varepsilon_{ijk} \hat{x}_k \quad (\text{formula})$$

((Proof))

$$\begin{aligned}
[\hat{L}_i, \hat{x}_j] &= \sum_{k,l} \epsilon_{ikl} [\hat{x}_k \hat{p}_l, \hat{x}_j] \\
&= \sum_{k,l} \epsilon_{ikl} (\hat{x}_k \hat{p}_l \hat{x}_j - \hat{x}_j \hat{x}_k \hat{p}_l) \\
&= \sum_{k,l} \epsilon_{ikl} (\hat{x}_k \hat{p}_l \hat{x}_j - \hat{x}_k \hat{x}_j \hat{p}_l) \\
&= \sum_{k,l} \epsilon_{ikl} \hat{x}_k [\hat{p}_l, \hat{x}_j] \\
&= -i\hbar \sum_{k,l} \epsilon_{ikl} \hat{x}_k \delta_{lj} \\
&= i\hbar \sum_{k,l} \epsilon_{ilk} \hat{x}_k \delta_{lj}
\end{aligned}$$

where ϵ_{ijk} is the Levi-Civita coefficient; $\epsilon_{ijk} = -\epsilon_{ikj}$. Thus, we have

$$[\hat{L}_i, \hat{x}_j] = i\hbar \sum_k \epsilon_{ijk} \hat{x}_k$$

Using this relation, we get

$$\begin{aligned}
[\hat{L}_x, \hat{x}] &= 0, & [\hat{L}_x, \hat{y}] &= i\hbar \hat{z}, & [\hat{L}_x, \hat{z}] &= -i\hbar \hat{y} \\
[\hat{L}_y, \hat{x}] &= -i\hbar \hat{z}, & [\hat{L}_y, \hat{y}] &= 0, & [\hat{L}_y, \hat{z}] &= i\hbar \hat{x} \\
[\hat{L}_z, \hat{x}] &= i\hbar \hat{z}, & [\hat{L}_z, \hat{y}] &= -i\hbar \hat{x}, & [\hat{L}_z, \hat{z}] &= 0
\end{aligned}$$

((Mathematica))

We show that the differential operators satisfy

$$[L_x, x] = 0, \quad [L_x, y] = i\hbar z, \quad [L_x, z] = -i\hbar y$$

Using Mathematica.

```
eq1 = Lx[x ψ[x, y, z] ] - x Lx[ψ[x, y, z] ] // Simplify
```

0

```
eq2 = Lx[y ψ[x, y, z] ] - y Lx[ψ[x, y, z] ] -  
i ħ z ψ[x, y, z] // FullSimplify
```

0

```
eq3 =  
Lx[z ψ[x, y, z] ] - z Lx[ψ[x, y, z] ] +  
i y ħ ψ[x, y, z] // Simplify
```

3 Commutation relation (II)

$$[\hat{L}^2, \hat{x}_k] = -i\hbar \sum_{i,j} \epsilon_{ijk} (\hat{L}_i \hat{x}_j + \hat{x}_j \hat{L}_i) \quad \text{(formula)}$$

((Proof))

$$\hat{L}^2 = \hat{L}_x^2 + \hat{L}_y^2 + \hat{L}_z^2$$

$$\begin{aligned} [\hat{L}_i^2, \hat{x}_j] &= \hat{L}_i \hat{L}_i \hat{x}_j - \hat{x}_j \hat{L}_i \hat{L}_i \\ &= \hat{L}_i (\hat{L}_i \hat{x}_j - \hat{x}_j \hat{L}_i) + (\hat{L}_i \hat{x}_j \hat{L}_i - \hat{x}_j \hat{L}_i \hat{L}_i) \\ &= \hat{L}_i [\hat{L}_i, \hat{x}_j] + [\hat{L}_i, \hat{x}_j] \hat{L}_i \\ &= i\hbar \sum_k \epsilon_{ijk} (\hat{L}_i \hat{x}_k + \hat{x}_k \hat{L}_i) \end{aligned}$$

or

$$[\hat{L}_1^2, \hat{x}_j] = i\hbar \sum_k \epsilon_{1jk} (\hat{L}_1 \hat{x}_k + \hat{x}_k \hat{L}_1)$$

$$[\hat{L}_2^2, \hat{x}_j] = i\hbar \sum_k \epsilon_{2jk} (\hat{L}_2 \hat{x}_k + \hat{x}_k \hat{L}_2)$$

$$[\hat{L}_3^2, \hat{x}_j] = i\hbar \sum_k \epsilon_{3jk} (\hat{L}_3 \hat{x}_k + \hat{x}_k \hat{L}_3)$$

Then we get

$$\begin{aligned}
[\hat{\mathbf{L}}^2, \hat{x}_k] &= [\hat{L}_1^2, \hat{x}_k] + [\hat{L}_2^2, \hat{x}_k] + [\hat{L}_3^2, \hat{x}_k] \\
&= i\hbar \sum_j \varepsilon_{1kj} (\hat{L}_1 \hat{x}_j + \hat{x}_j \hat{L}_1) + i\hbar \sum_k \varepsilon_{2kj} (\hat{L}_2 \hat{x}_j + \hat{x}_j \hat{L}_2) \\
&\quad + i\hbar \sum_j \varepsilon_{3kj} (\hat{L}_3 \hat{x}_j + \hat{x}_j \hat{L}_3) \\
&= i\hbar \sum_{i,j} \varepsilon_{ikj} (\hat{L}_i \hat{x}_j + \hat{x}_j \hat{L}_i) \\
&= -i\hbar \sum_{i,j} \varepsilon_{ijk} (\hat{L}_i \hat{x}_j + \hat{x}_j \hat{L}_i)
\end{aligned}$$

or

$$[\hat{\mathbf{L}}^2, \hat{x}_k] = -i\hbar \sum_{i,j} \varepsilon_{ijk} (\hat{L}_i \hat{x}_j + \hat{x}_j \hat{L}_i)$$

We have the following relations with differential operators.

$$\begin{aligned}
[\mathbf{L}^2, x] &= -i\hbar \varepsilon_{231} (L_2 x_3 + x_3 L_2) - i\hbar \varepsilon_{321} (L_3 x_2 + x_2 L_3) \\
&= -i\hbar (L_y z + z L_y) + i\hbar (L_z y + y L_z)
\end{aligned}$$

$$[\mathbf{L}^2, y] = -i\hbar (L_z x + x L_z) + i\hbar (L_x z + z L_x)$$

$$[\mathbf{L}^2, z] = -i\hbar (L_x y + y L_x) + i\hbar (L_y x + x L_y)$$

((**Mathematica**))

```

 $\text{eq1} = \text{Lsq}[x \psi[x, y, z]] - x \text{Lsq}[\psi[x, y, z]] +$ 
 $\frac{i \hbar}{2} (\text{Ly}[z \psi[x, y, z]] + z \text{Ly}[\psi[x, y, z]]) -$ 
 $\frac{i \hbar}{2} (\text{Lz}[y \psi[x, y, z]] + y \text{Lz}[\psi[x, y, z]]) // \text{FullSimplify}$ 
 $0$ 

 $\text{eq2} = \text{Lsq}[y \psi[x, y, z]] - y \text{Lsq}[\psi[x, y, z]] +$ 
 $\frac{i \hbar}{2} (\text{Lz}[x \psi[x, y, z]] + x \text{Lz}[\psi[x, y, z]]) -$ 
 $\frac{i \hbar}{2} (\text{Lx}[z \psi[x, y, z]] + z \text{Lx}[\psi[x, y, z]]) // \text{FullSimplify}$ 
 $0$ 

 $\text{eq3} = \text{Lsq}[z \psi[x, y, z]] - z \text{Lsq}[\psi[x, y, z]] +$ 
 $\frac{i \hbar}{2} (\text{Lx}[y \psi[x, y, z]] + y \text{Lx}[\psi[x, y, z]]) -$ 
 $\frac{i \hbar}{2} (\text{Ly}[x \psi[x, y, z]] + x \text{Ly}[\psi[x, y, z]]) //$ 
 $\text{FullSimplify}$ 
 $0$ 

```

4. Commutation relation (III)

$$[\hat{\mathbf{L}}^2, [\hat{\mathbf{L}}^2, \hat{x}_k]] = 2\hbar^2 (\hat{\mathbf{L}}^2 \hat{x}_k + \hat{x}_k \hat{\mathbf{L}}^2) \quad (\text{formula})$$

((Proof))

$$[\hat{\mathbf{L}}^2, [\hat{\mathbf{L}}^2, \hat{x}_k]] = -i\hbar \sum_{i,j} \epsilon_{ijk} [\hat{\mathbf{L}}^2, (\hat{L}_i \hat{x}_j + \hat{x}_j \hat{L}_i)]$$

with

$$\begin{aligned} [\hat{\mathbf{L}}^2, \hat{L}_i \hat{x}_j] &= \hat{\mathbf{L}}^2 \hat{L}_i \hat{x}_j - \hat{L}_i \hat{x}_j \hat{\mathbf{L}}^2 \\ &= \hat{L}_i \hat{\mathbf{L}}^2 \hat{x}_j - \hat{L}_i \hat{x}_j \hat{\mathbf{L}}^2 \\ &= \hat{L}_i [\hat{\mathbf{L}}^2, \hat{x}_j] \\ &= (-i\hbar) \hat{L}_i \sum_{l,m} \epsilon_{lmj} (\hat{L}_l \hat{x}_m + \hat{x}_m \hat{L}_l) \\ &= (-i\hbar) \sum_{l,m} \epsilon_{lmj} (\hat{L}_i \hat{L}_l \hat{x}_m + \hat{L}_i \hat{x}_m \hat{L}_l) \end{aligned}$$

and

$$\begin{aligned}
[\hat{\mathbf{L}}^2, \hat{x}_j \hat{L}_i] &= \hat{\mathbf{L}}^2 \hat{x}_j \hat{L}_i - \hat{x}_j \hat{L}_i \hat{\mathbf{L}}^2 \\
&= \hat{\mathbf{L}}^2 \hat{x}_j \hat{L}_i - \hat{x}_j \hat{\mathbf{L}}^2 \hat{L}_i \\
&= [\hat{\mathbf{L}}^2, \hat{x}_j] \hat{L}_i \\
&= (-i\hbar) \sum_{l,m} \varepsilon_{lmj} (\hat{L}_l \hat{x}_m \hat{L}_i + \hat{x}_m \hat{L}_l \hat{L}_i)
\end{aligned}$$

Using the above relations, we get

$$\begin{aligned}
[\hat{\mathbf{L}}^2, [\hat{\mathbf{L}}^2, \hat{x}_k]] &= (-i\hbar)^2 \sum_{i,j,l,m} \varepsilon_{ijk} \varepsilon_{lmj} (\hat{L}_i \hat{L}_l \hat{x}_m + \hat{L}_i \hat{x}_m \hat{L}_l + \hat{L}_l \hat{x}_m \hat{L}_i + \hat{x}_m \hat{L}_l \hat{L}_i) \\
&= \hbar^2 \sum_{i,j,l,m} \varepsilon_{ikj} \varepsilon_{lmj} (\hat{L}_i \hat{L}_l \hat{x}_m + \hat{L}_i \hat{x}_m \hat{L}_l + \hat{L}_l \hat{x}_m \hat{L}_i + \hat{x}_m \hat{L}_l \hat{L}_i)
\end{aligned}$$

with

$$\varepsilon_{ijk} \varepsilon_{lmj} = -\varepsilon_{jik} \varepsilon_{jlm}$$

Using the formula for the Levi-Civita

$$\sum_j \varepsilon_{jik} \varepsilon_{jlm} = \delta_{il} \delta_{km} - \delta_{im} \delta_{kl}$$

we have

$$\begin{aligned}
[\hat{\mathbf{L}}^2, [\hat{\mathbf{L}}^2, \hat{x}_k]] &= \hbar^2 \sum_{i,l,m} (\delta_{il} \delta_{km} - \delta_{im} \delta_{kl}) (\hat{L}_i \hat{L}_l \hat{x}_m + \hat{L}_i \hat{x}_m \hat{L}_l + \hat{L}_l \hat{x}_m \hat{L}_i + \hat{x}_m \hat{L}_l \hat{L}_i) \\
&= \hbar^2 \sum_i [(\hat{L}_i \hat{L}_i \hat{x}_k + \hat{x}_k \hat{L}_i \hat{L}_i + \hat{L}_i \hat{x}_k \hat{L}_i + \hat{L}_i \hat{x}_k \hat{L}_i) - (\hat{L}_i \hat{L}_k \hat{x}_i + \hat{L}_i \hat{x}_i \hat{L}_k + \hat{L}_k \hat{x}_i \hat{L}_i + \hat{x}_i \hat{L}_k \hat{L}_i)] \\
&= \hbar^2 (\hat{\mathbf{L}}^2 \hat{x}_k + \hat{x}_k \hat{\mathbf{L}}^2) + \hbar^2 \sum_{i,i} (\hat{L}_i \hat{x}_k \hat{L}_i + \hat{L}_i \hat{x}_k \hat{L}_i) - \hbar^2 \sum_i (\hat{L}_i \hat{L}_k \hat{x}_i + \hat{L}_i \hat{x}_i \hat{L}_k + \hat{L}_k \hat{x}_i \hat{L}_i + \hat{x}_i \hat{L}_k \hat{L}_i) \\
&= 2\hbar^2 (\hat{\mathbf{L}}^2 \hat{x}_k + \hat{x}_k \hat{\mathbf{L}}^2) - i\hbar^3 \sum_{i,l} \varepsilon_{ikl} [\hat{L}_i, \hat{x}_l] - \hbar^2 \sum_i (\hat{L}_i \hat{L}_k \hat{x}_i + \hat{L}_i \hat{x}_i \hat{L}_k + \hat{L}_k \hat{x}_i \hat{L}_i + \hat{x}_i \hat{L}_k \hat{L}_i) \\
&= 2\hbar^2 (\hat{\mathbf{L}}^2 \hat{x}_k + \hat{x}_k \hat{\mathbf{L}}^2)
\end{aligned}$$

(a)

$$[\hat{L}_i, \hat{x}_k] = i\hbar \sum_l \varepsilon_{ikl} \hat{x}_l, \quad [\hat{L}_k, \hat{x}_i] = i\hbar \sum_l \varepsilon_{kil} \hat{x}_l$$

or

$$\hat{L}_i \hat{x}_k - \hat{x}_k \hat{L}_i = i\hbar \sum_l \varepsilon_{ikl} \hat{x}_l$$

(b)

$$\begin{aligned} \hat{L}_i \hat{x}_k \hat{L}_i + \hat{L}_i \hat{x}_k \hat{L}_i &= \hat{L}_i [\hat{L}_i \hat{x}_k - i\hbar \sum_l \varepsilon_{ikl} \hat{x}_l] + [\hat{x}_k \hat{L}_i + i\hbar \sum_l \varepsilon_{ikl} \hat{x}_l] \hat{L}_i \\ &= \hat{L}_i \hat{L}_i \hat{x}_k + \hat{x}_k \hat{L}_i \hat{L}_i + \hbar^2 \sum_{l,m} \varepsilon_{ikl} \varepsilon_{ilm} \hat{x}_m \end{aligned}$$

(d)

$$\begin{aligned} \sum_i (\hat{L}_i \hat{x}_k \hat{L}_i + \hat{L}_i \hat{x}_k \hat{L}_i) &= \sum_i \{\hat{L}_i \hat{L}_i \hat{x}_k + \hat{x}_k \hat{L}_i \hat{L}_i - i\hbar \sum_{i,l} \varepsilon_{ikl} [\hat{L}_i, \hat{x}_l]\} \\ &= \hat{\mathbf{L}}^2 \hat{x}_k + \hat{x}_k \hat{\mathbf{L}}^2 + \hbar^2 \sum_{i,l,m} \varepsilon_{ikl} \varepsilon_{ilm} \hat{x}_m \end{aligned}$$

(d)

$$\hat{\mathbf{L}} \cdot \hat{\mathbf{r}} = \sum_i \hat{L}_i \hat{x}_i = 0, \quad \hat{\mathbf{r}} \cdot \hat{\mathbf{L}} = \sum_i \hat{x}_i \hat{L}_i = 0$$

So we need to show that

$$-i\hbar^3 \sum_{i,l} \varepsilon_{ikl} [\hat{L}_i, \hat{x}_l] - \hbar^2 \sum_i (\hat{L}_i \hat{L}_k \hat{x}_i + \hat{L}_i \hat{x}_i \hat{L}_k + \hat{L}_k \hat{x}_i \hat{L}_i + \hat{x}_i \hat{L}_k \hat{L}_i) = 0$$

where

$$\sum_{i,l} \varepsilon_{ikl} [\hat{L}_i, \hat{x}_l] = -2i\hbar x_k, \quad -i\hbar^3 \sum_{i,l} \varepsilon_{ikl} [\hat{L}_i, \hat{x}_l] = -i\hbar^3 (-2i\hbar x_k) = -2\hbar^4 x_k$$

$$-\hbar^2 \sum_i (\hat{L}_i \hat{L}_k \hat{x}_i + \hat{x}_i \hat{L}_k \hat{L}_i) = 2\hbar^4 x_k$$

$$-\hbar^2 \sum_i \hat{L}_i \hat{x}_i \hat{L}_k = 0 \quad -\hbar^2 \sum_i \hat{L}_k \hat{x}_i \hat{L}_i = 0$$

((Proof))

$$\hat{L}_i \hat{L}_k \hat{x}_i + \hat{x}_i \hat{L}_k \hat{L}_i = \hat{L}_i (\hat{x}_i \hat{L}_k + i\hbar \sum_l \varepsilon_{kil} \hat{x}_l) + (\hat{L}_k \hat{x}_i - i\hbar \sum_l \varepsilon_{kil} \hat{x}_l) \hat{L}_i$$

$$\begin{aligned} \sum_i (\hat{L}_i \hat{L}_k \hat{x}_i + \hat{x}_i \hat{L}_k \hat{L}_i) &= \sum_i [\hat{L}_i (\hat{x}_i \hat{L}_k + i\hbar \sum_l \varepsilon_{kil} \hat{x}_l) + (\hat{L}_k \hat{x}_i - i\hbar \sum_l \varepsilon_{kil} \hat{x}_l) \hat{L}_i] \\ &= \sum_i (\hat{L}_i \hat{x}_i \hat{L}_k + \hat{L}_k \hat{x}_i \hat{L}_i) + i\hbar \sum_{i,l} \varepsilon_{kil} [\hat{L}_i, \hat{x}_l] \\ &= i\hbar \sum_{i,l} \varepsilon_{kil} [\hat{L}_i, \hat{x}_l] \\ &= -2\hbar^2 \hat{x}_k \end{aligned}$$

((Mathematica))

We have the following relations with differential operators.

$$[\mathbf{L}^2, [\mathbf{L}^2, x]] = 2\hbar^2 (\mathbf{L}^2 x + x \mathbf{L}^2)$$

```
eq1 = Lsq[Lsq[x \psi[x, y, z]]] - Lsq[x Lsq[\psi[x, y, z]]] -
Lsq[x Lsq[\psi[x, y, z]]] + x Lsq[Lsq[\psi[x, y, z]]]] //
FullSimplify
```

$4 \hbar^4$

$$(x \psi[x, y, z] + 3 x z \psi^{(0, 0, 1)}[x, y, z] - x^3 \psi^{(0, 0, 2)}[x, y, z] - x y^2 \psi^{(0, 0, 2)}[x, y, z] + 3 x y \psi^{(0, 1, 0)}[x, y, z] + 2 x y z \psi^{(0, 1, 1)}[x, y, z] - x^3 \psi^{(0, 2, 0)}[x, y, z] - x z^2 \psi^{(0, 2, 0)}[x, y, z] + 2 x^2 \psi^{(1, 0, 0)}[x, y, z] - y^2 \psi^{(1, 0, 0)}[x, y, z] - z^2 \psi^{(1, 0, 0)}[x, y, z] + 2 x^2 z \psi^{(1, 0, 1)}[x, y, z] + 2 x^2 y \psi^{(1, 1, 0)}[x, y, z] - x (y^2 + z^2) \psi^{(2, 0, 0)}[x, y, z])$$

```

eq2 = 2  $\hbar^2$  (x Lsq[ψ[x, y, z]] + Lsq[x ψ[x, y, z]]) //  

FullSimplify

4  $\hbar^4$   

( x ψ[x, y, z] + 3 x z ψ^{(0,0,1)}[x, y, z] - x^3 ψ^{(0,0,2)}[x, y, z] -  

x y^2 ψ^{(0,0,2)}[x, y, z] + 3 x y ψ^{(0,1,0)}[x, y, z] +  

2 x y z ψ^{(0,1,1)}[x, y, z] - x^3 ψ^{(0,2,0)}[x, y, z] -  

x z^2 ψ^{(0,2,0)}[x, y, z] + 2 x^2 ψ^{(1,0,0)}[x, y, z] -  

y^2 ψ^{(1,0,0)}[x, y, z] - z^2 ψ^{(1,0,0)}[x, y, z] +  

2 x^2 z ψ^{(1,0,1)}[x, y, z] + 2 x^2 y ψ^{(1,1,0)}[x, y, z] -  

x (y^2 + z^2) ψ^{(2,0,0)}[x, y, z] )

```

eq1 - eq2 // Simplify

0

5. Commutators

We calculate the differential operators in the Cartesian coordinates

$$[L_z, r^2] = 0,$$

$$[L_z, \mathbf{p}^2] = 0,$$

$$[\mathbf{L}^2, \mathbf{p}^2] = 0$$

Using Mathematica.

((Mathematica))

```

eq1 =
Lz[ (x2 + y2 + z2) ψ[x, y, z] ] -
(x2 + y2 + z2) Lz[ ψ[x, y, z] ] // Simplify

0

eq2 = Lz[Peq[ψ[x, y, z]]] - Peq[Lz[ψ[x, y, z]]] //
Simplify

0

eq3 =
Leq[Peq[ψ[x, y, z]]] - Peq[Leq[ψ[x, y, z]]] // Simplify

0

```

REFERENCES

- A. Goswami, Quantum Mechanics, 2nd edition (Wm. C. Brown Publishers, 1997).
J. Binney and D. Skinner, The Physics of Quantum Mechanics (Oxford, 2014).

APPENDIX

```

Clear["Global`"];
ux = {1, 0, 0};
uy = {0, 1, 0};
uz = {0, 0, 1};
r = {x, y, z};
R = Sqrt[r.r];
L := (-I \hbar (Cross[r, Grad[#, {x, y, z}], "Cartesian"]]) & // Simplify;
p = (-I \hbar Grad[#, {x, y, z}, "Cartesian"] &) // Simplify;
Lx := (ux.L[#]) &;
Ly := (uy.L[#]) &;
Lz := (uz.L[#]) &;
px := (ux.p[#]) &;
py := (uy.p[#]) &;
pz := (uz.p[#]) &;
Lap := Laplacian[#, {x, y, z}, "Cartesian"] &;
Gra := Grad[#, {x, y, z}, "Cartesian"] &;
Cur := Curl[#, {x, y, z}, "Cartesian"] &;
Diva := Div[#, {x, y, z}, "Cartesian"] &;
prq := (-I \hbar) \frac{1}{2} \left( \frac{r}{R}.Gra[#] + Diva\left[\frac{r}{R}\# \right] \right) &;

```

```
Lsq := (Nest[Lx, #, 2] + Nest[Ly, #, 2] + Nest[Lz, #, 2]) &;
```

Goswami Quantum mechanics

Problem A-1

```
s11 = Lsq[xψ[x, y, z]] - x Lsq[ψ[x, y, z]] // Simplify
```

$$2 \hbar^2 (x \psi[x, y, z] + x z \psi^{(0,0,1)}[x, y, z] + x y \psi^{(0,1,0)}[x, y, z] - y^2 \psi^{(1,0,0)}[x, y, z] - z^2 \psi^{(1,0,0)}[x, y, z])$$

```
s12 = Lsq[yψ[x, y, z]] - y Lsq[ψ[x, y, z]] // Simplify
```

$$2 \hbar^2 (y \psi[x, y, z] + y z \psi^{(0,0,1)}[x, y, z] - x^2 \psi^{(0,1,0)}[x, y, z] - z^2 \psi^{(0,1,0)}[x, y, z] + x y \psi^{(1,0,0)}[x, y, z])$$

```
s13 = Lsq[zψ[x, y, z]] - z Lsq[ψ[x, y, z]] // Simplify
```

$$2 \hbar^2 (z \psi[x, y, z] - (x^2 + y^2) \psi^{(0,0,1)}[x, y, z] + z (y \psi^{(0,1,0)}[x, y, z] + x \psi^{(1,0,0)}[x, y, z]))$$

```
Lsq := (Nest[Lx, #, 2] + Nest[Ly, #, 2] + Nest[Lz, #, 2]) &;
```

Goswami Quantum mechanics

Problem A-1

```
s11 = Lsq[xψ[x, y, z]] - x Lsq[ψ[x, y, z]] // Simplify
```

$$2 \hbar^2 (x \psi[x, y, z] + x z \psi^{(0,0,1)}[x, y, z] + x y \psi^{(0,1,0)}[x, y, z] - y^2 \psi^{(1,0,0)}[x, y, z] - z^2 \psi^{(1,0,0)}[x, y, z])$$

```
s12 = Lsq[yψ[x, y, z]] - y Lsq[ψ[x, y, z]] // Simplify
```

$$2 \hbar^2 (y \psi[x, y, z] + y z \psi^{(0,0,1)}[x, y, z] - x^2 \psi^{(0,1,0)}[x, y, z] - z^2 \psi^{(0,1,0)}[x, y, z] + x y \psi^{(1,0,0)}[x, y, z])$$

```
s13 = Lsq[zψ[x, y, z]] - z Lsq[ψ[x, y, z]] // Simplify
```

$$2 \hbar^2 (z \psi[x, y, z] - (x^2 + y^2) \psi^{(0,0,1)}[x, y, z] + z (y \psi^{(0,1,0)}[x, y, z] + x \psi^{(1,0,0)}[x, y, z]))$$

Goswami Quantum mechanics

Problem A - 2

```
s21 = Lx[ψ[x, y, z]] - (y pz[ψ[x, y, z]] - z py[ψ[x, y, z]]) // Simplify
```

0

```
s22 = Ly[ψ[x, y, z]] - (z px[ψ[x, y, z]] - x pz[ψ[x, y, z]]) // Simplify
```

0

```
s23 = Lz[ψ[x, y, z]] - (x py[ψ[x, y, z]] - y px[ψ[x, y, z]]) // Simplify
```

0

Goswami Quantum mechanics

Problem A-3

s31 = Lsq[xψ[x, y, z] - x Lsq[ψ[x, y, z]] // Simplify

$$2 \hbar^2 (x \psi[x, y, z] + x z \psi^{(0,0,1)}[x, y, z] + x y \psi^{(0,1,0)}[x, y, z] - y^2 \psi^{(1,0,0)}[x, y, z] - z^2 \psi^{(1,0,0)}[x, y, z])$$

s32 =

-i ħ (Ly[zψ[x, y, z]] + z Ly[ψ[x, y, z]]) +
i ħ (Lz[yψ[x, y, z]] + y Lz[ψ[x, y, z]]) // Simplify

$$2 \hbar^2 (x \psi[x, y, z] + x z \psi^{(0,0,1)}[x, y, z] + x y \psi^{(0,1,0)}[x, y, z] - y^2 \psi^{(1,0,0)}[x, y, z] - z^2 \psi^{(1,0,0)}[x, y, z])$$

s32 - s31 // Simplify

0

Goswami Quantum mechanics

Problem A - 4

```
s41 = Lsq[Lsq[xψ[x, y, z]]] - Lsq[x Lsq[ψ[x, y, z]]] - Lsq[x Lsq[ψ[x, y, z]]] +  
x Lsq[Lsq[ψ[x, y, z]]] // FullSimplify
```

$$4\hbar^4 \left(x\psi[x, y, z] + 3xz\psi^{(0,0,1)}[x, y, z] - x^3\psi^{(0,0,2)}[x, y, z] - \right. \\ \left. xy^2\psi^{(0,0,2)}[x, y, z] + 3xy\psi^{(0,1,0)}[x, y, z] + 2xyz\psi^{(0,1,1)}[x, y, z] - \right. \\ \left. x^3\psi^{(0,2,0)}[x, y, z] - xz^2\psi^{(0,2,0)}[x, y, z] + 2x^2\psi^{(1,0,0)}[x, y, z] - \right. \\ \left. y^2\psi^{(1,0,0)}[x, y, z] - z^2\psi^{(1,0,0)}[x, y, z] + 2x^2z\psi^{(1,0,1)}[x, y, z] + \right. \\ \left. 2x^2y\psi^{(1,1,0)}[x, y, z] - x(y^2 + z^2)\psi^{(2,0,0)}[x, y, z] \right)$$

```
s42 = 2h^2 (x Lsq[ψ[x, y, z]] + Lsq[xψ[x, y, z]]) // FullSimplify
```

$$4\hbar^4 \left(x\psi[x, y, z] + 3xz\psi^{(0,0,1)}[x, y, z] - x^3\psi^{(0,0,2)}[x, y, z] - \right. \\ \left. xy^2\psi^{(0,0,2)}[x, y, z] + 3xy\psi^{(0,1,0)}[x, y, z] + 2xyz\psi^{(0,1,1)}[x, y, z] - \right. \\ \left. x^3\psi^{(0,2,0)}[x, y, z] - xz^2\psi^{(0,2,0)}[x, y, z] + 2x^2\psi^{(1,0,0)}[x, y, z] - \right. \\ \left. y^2\psi^{(1,0,0)}[x, y, z] - z^2\psi^{(1,0,0)}[x, y, z] + 2x^2z\psi^{(1,0,1)}[x, y, z] + \right. \\ \left. 2x^2y\psi^{(1,1,0)}[x, y, z] - x(y^2 + z^2)\psi^{(2,0,0)}[x, y, z] \right)$$

```
s42 - s41 // Simplify
```

```
0
```

Show that s51+s52+s53+s54+s55==0

```
s51 =
-i  $\hbar^3$  (-Ly[z  $\psi[x, y, z]$ ] + z Ly[ $\psi[x, y, z]$ ] + Lz[y  $\psi[x, y, z]$ ] -
y Lz[ $\psi[x, y, z]$ ]) // Simplify

-2 x  $\hbar^4$   $\psi[x, y, z]$ 

s52 = - $\hbar^2$  (Lx[Lx[x  $\psi[x, y, z]$ ]] + Ly[Lx[y  $\psi[x, y, z]$ ]] + Lz[Lx[z  $\psi[x, y, z]$ ]]) // FullSimplify

 $\hbar^4$  (2 x  $\psi[x, y, z]$  + x z  $\psi^{(0,0,1)}[x, y, z]$  +
x y  $\psi^{(0,1,0)}[x, y, z]$  - (y2 + z2)  $\psi^{(1,0,0)}[x, y, z]$ )

s53 = - $\hbar^2$  (Lx[x Lx[ $\psi[x, y, z]$ ]] + Ly[y Lx[ $\psi[x, y, z]$ ]] + Lz[z Lx[ $\psi[x, y, z]$ ]]) // FullSimplify

0

s54 = - $\hbar^2$  (Lx[x Lx[ $\psi[x, y, z]$ ]] + Lx[y Ly[ $\psi[x, y, z]$ ]] + Lx[z Lz[ $\psi[x, y, z]$ ]]) // FullSimplify

0

s54 = - $\hbar^2$  (Lx[x Lx[ $\psi[x, y, z]$ ]] + Lx[y Ly[ $\psi[x, y, z]$ ]] + Lx[z Lz[ $\psi[x, y, z]$ ]]) // FullSimplify

0

s55 = - $\hbar^2$  (x Lx[Lx[ $\psi[x, y, z]$ ]] + y Lx[Ly[ $\psi[x, y, z]$ ]] + z Lx[Lz[ $\psi[x, y, z]$ ]]) // FullSimplify

 $\hbar^4$  (-x (z  $\psi^{(0,0,1)}[x, y, z]$  + y  $\psi^{(0,1,0)}[x, y, z]$ ) + (y2 + z2)  $\psi^{(1,0,0)}[x, y, z]$ )

s51 + s52 + s53 + s54 + s55 // FullSimplify

0

s52 + s55 // Simplify

2 x  $\hbar^4$   $\psi[x, y, z]$ 

Lx[x  $\psi[x, y, z]$ ] - x Lx[ $\psi[x, y, z]$ ] // Simplify

0
```